Resonance Raman Spectrum of Dibromo^[N.N'-bis-1] **(salicylidene)propylenediaminatooxovanadium(V)] copper(I)*Methanol**

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Recently, Okawa and Kida [I] suggested that the strong electronic absorption band at 17.9 \times 10³ cm^{-1} of dibromo [N,N'-bis(salicylidene)propylenediaminatooxovanadium (V)] copper (I) \cdot methanol

in methanol solution may be due to the charge-transfer from the $p\pi$ -orbital of the phenolic oxygen to the lowest vacant orbital (d_{xy}) of the vanadium(V) atom. In the present communication, we wish to report the resonance Raman spectrum of this compound which provides direct evidence to support their assignment.

Experimental

The binuclear complex, $(H, H, Me)VOCuBr₂$ -CH,OH, was prepared by the method reported previously [l] . Both electronic and resonance Raman spectra of this complex were measured as KBr pellets. The electronic spectrum was recorded on a Gary Model 14 spectrophotometer. The resonance Raman spectra were measured by using a Spex Model 1401 double monochromator. Detection was made with a cooled RCA C31034A photo-multiplier tube with DC amplification. Excitations at 454.5 and 514.5 nm were made by a Spectra-Physics Model 164 Argon-ion laser and that at 578.8 nm was made by a Spectra-Physics Model 365 CW dye-laser (Rhodamine 6G) pumped by the above 4W Argon-ion laser. To compare relative intensity, the internal standard $(NaNO₃)$ was mixed homogeneously with the sample in a KBr pellet. The rotating sample technique was used to avoid thermal decomposition of the complex. The infrared spectra were measured on a Beckman IR 12 infrared spectrophotometer. The spectra of the complexes were measured as KBr pellets while that of the pure ligand was obtained as a neat liquid.

 \mathcal{C} is a known per \mathcal{C} . OH in a KBr pellet.

 $\sum_{i=1}^{\infty}$ CSOHance Nan

Results and Discussion

Figure 1 illustrates the electronic spectrum of (H, H, Me) VOCuBr₂ \cdot CH₃OH in a KBr pellet. Relative to the spectrum in solution [I], the bands observed in the KBr pellet are broader and flatter and the absorption maxima are shifted slightly. Even so, it is clear that the bands at 16×10^3 and 29×10^3 cm^{-1} in the solid state correspond to those observed at 17.9 \times 10³ and 27 \times 10³ cm⁻¹, respectively, in methanol solution. Figure 2 shows the resonance Raman spectra of the same compound in a KBr pellet with the laser excitations at 454.5, 514.5 and 578.8 nm. The bands marked by asterisks (1388, 100 and 726 cm⁻¹) are due to the internal standard, \sim \sim $NaNO₃$.
At the 454.5 nm excitation, the compound

exhibits a number of strong bands as is shown in the $\frac{1}{2}$ in trace of $\sum_{i=1}^{n}$ $\sum_{i=1}^{n}$ is safe to assign the strongest band at 977 cm^{-1} to the $V=0$ stretching and the bands between 1700 and 1000 cm⁻¹ to and all the bands between 1700 and 1000 cm^{-1} to the vibrations of the Schiff base ligand [2]. The

strong band at 647 cm^{-1} with a shoulder at 662 cm^{-1} is assigned to the V-O (phenolic) stretching mode. This assignment is supported by the observation that its infrared spectrum exhibits a band at 660 cm^{-1} with a shoulder at 648 cm^{-1} , both of which are absent in the infrared spectrum of (H, H, Me)Cu. The small splitting of this mode into two bands may be attributed to the crystal lattice effect. These frequencies are much higher than the V-O stretching frequencies of $VO(acac)$, (480 cm^{-1}) [3] due to the differences in the oxidation state of the vanadium atom $(+V \nu s. +IV)$ and in the nature of the oxygen atom (phenolic YS. ketonic). In fact, the antisymmetric and symmetric VOV stretching frequencies of the $[V_2O_7]^{4-}$ ion in which the vanadium is in the +V state are $780-680$ and $545-500$ cm⁻¹, respective**lY 141.**

At the 514.5 nm excitation (middle trace), the V=O stretching band at 977 cm $^{-1}$ becomes weaker together with the ligand bands between 1700 and 1000 cm^{-1} . On the other hand, the V-O (phenolic) stretching band at 647 cm⁻¹ becomes stronger relative to other bands. Also, new bands emerge at 576 and 490 cm^{-1} which correspond to the infrared bands at 577 and 491 cm^{-1} . Since these bands are not present in the infrared spectrum of (H, H, Me)Cu, they are probably due to other $VO₃N₂$ L skeletal modes such as the V-N and V-L stretching modes.

When the exciting frequency (578.8 nm) approaches the absorption maximum of the lower energy electronic band at 16×10^3 cm⁻¹, the V=O stretching band at 977 cm^{-1} is almost extinct. So are some ligand vibrations between 1700 and 1000 cm^{-1} . On the other hand, the V-0 (phenolic) stretching band at 647 cm^{-1} becomes almost three times stronger than the internal standard band at 1069

 cm^{-1} (top trace). The other skeletal bands at 576 and 490 cm^{-1} are also enhanced.

Previously, Okawa and Kida [l] assigned the lower energy electronic band of (H, H, Me)VOCu-Br₂ CH₃OH to the V-O (phenolic) charge-transfer transition based on the following observations: (1) the intensity of this band (log $\epsilon \approx 3$) is unusually high for a d-d transition, (2) the vanadium(V) complex of N,N'-bis(salicylidene)ethylenediamine also exhibits a strong band at 17.9×10^3 cm⁻¹, and (3) such bands are not observed for analogous comppounds in which the VO group is replaced by other metals.

According to the theories of resonance Raman scattering, only those vibrations which are localized on a particular chromophoric group are enhanced when the energy of the exciting line approaches that of the electronic transition which originates in this chromophore (selectivity) $[5, 6]$. Then the dramatic changes in the resonance Raman spectrum of (H, H, Me)VOCuBr₂ \cdot CH₃OH described above can only be understood if we assign the lower energy band to the V-O (phenolic) charge-transfer transition as previously suggested.

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